

## Exercise-sheet 1 (May 9, 2017)

### 1 In-class exercises. Free phonons: a short tour

In this first session we will revisit the classical vibrations of a systems of ions within the harmonic approximation, and their quantization into phonons. As an example, we will consider the *jellium model* and determine its normal modes.

### 2 Homework - due date: May 16, 2017 (35 points).

#### 2.1 Revisiting Debye's model of phonons (10 points).

Consider an elastic medium of volume  $V$ . The frequency associated to the  $3N$  normal modes of a  $N$ -atom solid can be approximated by the lowest  $3N$  normal frequencies of the elastic medium.

- (a) Find the number of normal modes  $g(\omega)d\omega$  for an elastic medium whose frequencies lie between  $\omega$  and  $d\omega$ , assuming the system has periodic boundary conditions.
- (b) In a uniform medium, the normal frequencies have no upper limit. In a real crystal however, the wavenumber propagation cannot exceed the reciprocal of the interparticle spacing. Hence, there exist a natural high frequency cutoff  $\omega_D$ . Find  $\omega_D$  from the condition

$$3N = \int_0^{\omega_D} g(\omega)d\omega, \quad (1)$$

and estimate the corresponding wavelength  $\lambda_D$ .

- (c) Show the total energy of the system, defined as  $E = \sum_i \hbar\omega_i \langle n_i \rangle$ , with  $\langle n_i \rangle$  the Bose function, can be written as

$$E = \alpha \int_0^{T_D/T} dx \frac{x^3}{e^x - 1}, \quad (2)$$

where  $\alpha = \alpha(T)$  is to be determined and  $kT_D \equiv \hbar\omega_D$ .

- (d) Find the leading terms of the total energy and specific heat in both, the high ( $T \gg T_D$ ) and low ( $T \ll T_D$ ) temperature expansion of Eq. (2). You may need to use  $\zeta(4) = \frac{\pi^4}{90}$ , with  $\zeta(n)$  the Riemann zeta function.

## 2.2 Free electron Green's function (10 points).

Consider the free electron Hamiltonian

$$\hat{H} = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma}, \quad \epsilon_{\mathbf{k}} = \frac{k^2}{2m}, \quad (3)$$

where  $\hat{c}_{\mathbf{k}\sigma}$  ( $\hat{c}_{\mathbf{k}\sigma}^\dagger$ ) annihilates (creates) a  $\sigma$ -spin electron in a plane wave state  $\mathbf{k}$ .

The single particle Green's function is defined as the ground state expectation value of the *time-ordered product*

$$G^{(0)}(\mathbf{k}t) = -i\langle 0|T\hat{c}_{\mathbf{k}}(t)\hat{c}_{\mathbf{k}}^\dagger(0)|0\rangle, \quad (4)$$

where  $|0\rangle$  labels the filled Fermi sea and  $\hat{c}_{\mathbf{k}}(t)$  is a Heisenberg operator, with  $\hat{c}_{\mathbf{k}}(t) = e^{i\hat{H}t}\hat{c}_{\mathbf{k}}e^{-i\hat{H}t}$ . Furthermore, the time ordering operator is defined as

$$T\hat{a}(t)\hat{b}(t') = \theta(t-t')\hat{a}(t)\hat{b}(t') \pm \theta(t'-t)\hat{b}(t')\hat{a}(t), \quad (5)$$

where the upper (lower) sign refers to bosons (fermions) and  $\theta(t)$  is the Heaviside step function.

(a) Use Heisenberg's equation of motion to find the explicit time dependence of  $\hat{c}_{\mathbf{k}}$ .

(b) Show

$$G^{(0)}(\mathbf{k}t) = -i\theta(t)(1-f_{\mathbf{k}})e^{-i\epsilon_{\mathbf{k}}t} + i\theta(-t)f_{\mathbf{k}}e^{-i\epsilon_{\mathbf{k}}t} \quad (6)$$

where  $f_{\mathbf{k}}$  is the Fermi function.

(c) Show the time Fourier transform of Eq. (6) is given by

$$G^{(0)}(\mathbf{k}\omega) = \frac{1}{\omega - \epsilon_{\mathbf{k}} + i\eta_{\mathbf{k}}}, \quad (7)$$

where  $\eta_{\mathbf{k}}$  is a positive (negative) infinitesimal for  $\mathbf{k}$  greater (less) than the Fermi momentum  $k_F$ .

## 2.3 Free phonon Green's function (10 points).

Consider the free phonon Hamiltonian

$$\hat{H} = \sum_{\mathbf{q}} \Omega_{\mathbf{q}} \hat{a}_{\mathbf{q}}^\dagger \hat{a}_{\mathbf{q}}, \quad (8)$$

where  $\hat{a}_{\mathbf{q}}$  ( $\hat{a}_{\mathbf{q}}^\dagger$ ) annihilates (creates) a phonon with wave vector  $\mathbf{q}$ .

The phonon Green's function is defined as the ground state expectation value of the *time-ordered product*

$$D^{(0)}(\mathbf{q}t) = -i\langle 0|T\hat{A}_{\mathbf{q}}(t)\hat{A}_{\mathbf{q}}^{\dagger}(0)|0\rangle, \quad (9)$$

where  $\hat{A}_{\mathbf{q}} = a_{\mathbf{q}} + a_{-\mathbf{q}}^{\dagger}$  and  $T$  is defined in Eq. (5). Note that the ground state for a system of free phonons is the empty state.

Show

$$D^{(0)}(\mathbf{q}\omega) = \frac{2\Omega_{\mathbf{q}}}{\omega^2 + \Omega_{\mathbf{q}}^2 + i\delta}. \quad (10)$$

## 2.4 Matrix order parameter (5 points).

Consider the matrix order parameter

$$\Delta = (d_0\sigma_0 + \mathbf{d} \cdot \boldsymbol{\sigma}) i\sigma^2, \quad (11)$$

where  $\sigma^0$  is the  $2 \times 2$  identity matrix,  $\sigma^i$  ( $i = 1, 2, 3$ ) the Pauli matrices, and  $d_j$  ( $j = 0, 1, 2, 3$ ) are complex numbers. Show the gauge-invariant matrix

$$\Delta\Delta^{\dagger} = (|d_0|^2 + |\mathbf{d}|^2)\sigma^0 + \mathbf{q} \cdot \boldsymbol{\sigma}, \quad (12)$$

where  $\mathbf{q}$  is to be determined.